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Numerical Simulation of Casson Micropolar Fluid Flow Over an Inclined Surface Through Porous Medium

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Abstract

Numerical Simulations of Casson Micropolar nanofluid flow over a slanted surface through porous medium is debated in this article. The governing Partial differential equations with various parameters modified to Ordinary differential equations. The set of equations is determined numerically by using MATLAB bvp4c method. The fallout of velocity, Temperature, Concentration and Micro polar on Magnetic, Porous medium profiles is exposed graphically. The implication of Friction drag, Nusselt, and Sherwood numerals for various Magnetic, Casson parameter and Porosity quantities is high lightened in the table.

Keywords: Chemical Reaction, Micro Polar, Nano Fluid, Numerical Analysis, Porous Medium

1.0 Introduction

The heat transmission beyond a porous sheet has attracted to its expanding applications in corporation and modern technologies. The MHD flow of mechanized carrying out, insoluble fluid a wide plate that is evenly accelerated and isolated has a precise solution are investigated by Bala Anki Reddy¹. Nonlinear heat transfer attribute on thermal radiation-influenced magneto hydrodynamic Casson fluid flow in a hall effect are scrutinized by Basavaraj et al., and Deivanayaki et al.,^{2,3}. The MHD Rayleigh problems with an angled electromagnetic field to determine the impression of a viscous dissipation are studied by Fazle et al.,⁴. The heating process and Brownian decay coefficients in the border absorption, the residual of partial slip on the velocity region, as well as turbulent heat barrier case are scrutinized by Hemamalini *et al.*,⁵. The fallout of pedesis diffusion and thermophoretic on dynamically conductive varied flow Casson fluid acquired by a moving wedge are probed by Oyelakin et al.,6. A arithmetic review on Casson

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nanofluid horizontal huge surface with a magnetic field effects discussed. The analysis takes into account slip and heat boundary cases are discussed by Imran et al.,7. The reaction of thermal production and dispersion on MHD in the occupation of heat radiation and synthetic attitude on Casson fluid flow over a perpendicular plate that is oscillating and implanted through porous substances are investigated by Kamran et al.,8. The Buongiorno nanofluid approach, which are examined the heated performance of the smooth flows under motion of Brownian and thermophoretic phenomena and the operation of nanomaterial are discussed by Hari et al., Anwar et al and Venkata Ramudu et al.,^{9,10,12}. The development of warmth pitch initiation of porous permeability iron deportation on sub heated regular spongy layer is measured by Thomas et al.,¹¹. The above studies did not take into account of the Numerical findings on casson micropolar nano fluid flow extinct an inclined surface through porous medium. To the best of their understanding, no prior study comprises concepts which are analogous to those in this work.

The innovation is occupying on the viscous diversion for a non-Newtonian Casson, and the effectiveness of synthetic processes carried on by nanoparticle migrations in the field of producing warmth micropolar nanofluid sample infused in an absorptive medium and exposed to a magnetic field, warmth radiation, and a change in velocity aspect.

2.0 Mathematical Analysis

The subsequent equations for the continuity, momentum, energy and angular momentum

$$\frac{\partial u}{\partial x} + \frac{\partial u}{\partial y} = 0$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = v \left(1 + \frac{1}{\beta} \right) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{\partial^2 u}{\partial y^2} \right) \cdot \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta}) \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1}{\beta} \left(\frac{k_T}{\rho} \right) \frac{\partial N}{\partial y} + (g \beta_T (T - T_\infty) + \frac{1$$

$$u\frac{\partial N}{\partial x} + v\frac{\partial N}{\partial y} = \frac{\gamma}{\rho} \frac{\partial^2 N}{\partial y^2} - \left(\frac{k_1}{\rho}\right) \left(2N + \frac{\partial u}{\partial y}\right)$$
(2)

$$\rho C_{p} \left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = k \frac{\partial^{2} T}{\partial y^{2}} - \frac{\partial q_{r}}{\partial y} + q'' + \nu \left(1 + \frac{1}{\beta} \right) \rho \left(\frac{\partial u}{\partial y} \right)^{2} + \frac{\rho D_{m} k_{T}}{C_{s}} \frac{\partial^{2} C}{\partial y^{2}}$$
(3)

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_m \frac{\partial^2 C}{\partial y^2} - k^* (C - C_\infty) + \frac{D_m k_T}{T_m} \frac{\partial^2 T}{\partial y^2}$$
(4)

The irregular heat parameter is defined by

$$q''' = \frac{ku_{w}}{xv} \left(A^{*} (T_{w} - T_{\infty}) f' + B^{*} (T - T_{\infty}) \right)$$
(5)

Here A^* and B^* serve heat generation and the internal heat absorption. The warmth discrepancy between T and stream warmth T_{∞} is tiny, and by removing components of higher order, Taylor's simplification concerning T_{∞} is $T_w = T_{\infty} + bx$

The equation is the result of using Rosland's similarity for thermal emission,

$$q_r = -\frac{16}{3} \frac{\sigma^*}{k^*} T^3 \frac{\partial T}{\partial y} \tag{6}$$

The confines condition as follows

$$u = u_{w} + L \frac{\partial u}{\partial y}, -k \frac{\partial T}{\partial y} = h_{1} (T_{w} - T), -N = -m \frac{\partial u}{\partial y}, D_{m} \frac{\partial C}{\partial y} = h_{2} (C_{w} - C), \text{ as } y = 0$$
$$u \to 0, T \to T_{w}, C \to C_{w} \text{ as } y = \infty$$
(7)

Where L signifes molecular free mean groove, h_1 and h_2 concentus of deportation and eccentric transmission, C_w liquid's solidification at its outermost layer.

Consider the similarity transforms,

$$\varphi = \sqrt{av}xf(\eta), u = \frac{\partial\varphi}{\partial y} = axf'(\eta), N = ax\left(\sqrt{\frac{a}{v}}\right)h(\eta),$$
$$v = -\frac{\partial\varphi}{\partial x} = \sqrt{av}f(\eta)$$
$$\theta(\eta) = \frac{T - T_{\infty}}{T_w - T_{\infty}}, \qquad \phi(\eta) = \frac{C - C_{\infty}}{C_w - C_{\infty}},$$
$$T = T_{\infty}(1 + (\theta_w - 1)\theta)$$
(8)

Using Equation of (8) the equations. (2)-(4) are reconstruct as

$$\left(1 + \frac{1}{\beta}\right)f''' + ff'' - (f')^2 + Kh' + (Gr_T\theta + Gr_c\phi)\cos\alpha - (M + k_p)f' = 0$$
(9)

$$\left(1 + \frac{K}{2}\right)h'' + fh' - f\dot{h} + K(2h + f'') = 0$$
(10)

$$[1 + N_r \{1 + 3(\theta_w - 1) + 3(\theta_w - 1)^2 \theta^2 + (\theta_w - 1)^3 \theta^3\}]\theta'' +$$

$$N_{r}[(\theta_{w} - 1)(\theta')^{2} + 6(\theta_{w} - 1)^{2}\theta(\theta')^{2} + 3(\theta_{w} - 1)^{3}\theta^{2}(\theta')^{2}] + P_{r}f\theta' - P_{r}f'\theta + A^{*}f' + B^{*}\theta + P_{r}Du\phi'' + P_{r}E_{c}\left(1 + \frac{1}{\beta}\right)(f'')^{2} = 0 - (11)$$

$$\phi'' + Scf \phi' - k_r Sc \phi + ScSr \theta'' = 0$$
(12)

Where

$$Gr_T = \frac{g\beta_T(T_w - T_\infty)}{a^2 x}, N_r = \frac{16\sigma^* T_\infty^3}{3kk^*},$$

$$Gr_{C} = \frac{g\beta_{C}(C_{w} - C_{\infty})}{a^{2}x}, Pr = \frac{v\rho C_{p}}{k} = \frac{\mu C_{p}}{k}, M = \frac{\sigma B_{0}}{a\rho}$$
$$Sr = \frac{D_{m}K_{T}}{vT_{m}} \left(\frac{T_{w} - T_{\infty}}{C_{w} - C_{\infty}}\right), Sc = \frac{v}{D_{m}},$$

$$Du = \frac{D_m K_T (C_w - C_\infty)}{C_s C_p \upsilon (T_w - T_\infty)}, E_c = \frac{u^2}{C_p (T_w - T_\infty)}, K_r = \frac{k^*}{a}$$
(13)

The changed confines conditions as follows

$$f' = 1 + \gamma f''$$
, $f = 0$, $\theta' = Bi_T(1 - \theta)$, $\phi' = Bi_c(1 - \phi)$ as
 $\eta \to 0$,

$$f'(\eta) = 0, \ \theta(\eta) = 0, \ \phi(\eta) = 0 \text{ as } \eta \to \infty$$
 (14)

The friction drag, Nusselt and Sherwood numeral are the physical confines need attention and defined¹² as follows

$$C_{f_x} R e_x^{\frac{1}{2}} = \left(1 + \frac{1}{\beta}\right) f''(0) \tag{15}$$

$$NuRe_{x}^{\frac{1}{2}} = -(1 + N_{r}(\theta_{w})^{3})\theta'(0)$$
⁽¹⁶⁾

$$Sh Re_x^{\overline{2}} = - \phi'(0) \tag{17}$$

Here, $Re_x = \frac{xu_w}{v}$ signifies the Reynolds Number.

3.0 Results and Discussion

The nonlinear equations (9) – (14) are disclosed numerically by MATLAB bvp4c method. Consider the numerical values for the graphical results shown below, $\beta = 1.5, Gr_T = 0.8, M = 1.5, Nr = 0.2,$ $A^* = B^* = 0.1, Pr = 7, \theta_w = 2, Gr_c = 0.8,$



Figure 1. Velocity on diverse of M.

BiT = 2, BiC = 2, Sc = 0.6, Kr = 0.2,

Sr = 0.2, Ec = 0.2 and Du = 0.1. The significance of different deep flow metaphors is demonstrated by friction with the skin visuals, Sherwood and Nusselt computation, and cosmetic regards phases. Figures 1 & 4 display consequences of the magnetic on the micropolar depiction and velocity. A drop in velocity and micro polarity is noticed as soon as the magnetic criterion (M values) is increased. Due to its significant on the movement of fluids, M generates a force as the Lorentz force when imposed through the middle in the preferred direction of flow. Hydrodynamics, plasma accelerators, MHD accelerators, and other technical domains all use the Lorentz force in the electromagnetic force. As consequently, as the magnetic coefficient grows, the



Figure 2. Temperature on diverse of M.



Figure 3. Concentration on diverse of M



Figure 4. Angular velocity on diverse of M.



Figure 5. Velocity on diverse of β .



Figure 6. Temperature on diverse of β .



Figure 7. Concentration on diverse of β .



Figure 8. Angular velocity on diverse of β .



Figure 9. Velocity on *diverse* of K_n

Lorentz force leads the velocity and boundary layer density to decrease. Figures 2 & 3 show when a magnetic parameter grows, the warmth and absorption rise. As the outcome, a force can be created on the tension gradient of the Nano fluid content causing it to subsequently absorb some heat energy.



Figure 10. Temperature on *diverse* of K_b.

Table 1. Values of C_{fx} , Nu and Sh

Figures 5-8 depict Casson parameter on the flow fields					
momentum, warmth, absorption and micro rotation.					
Velocity profile and angular velocity declined for larger					
β while θ , \emptyset are embellished for the raise β . Figures 9-12					
display the velocity profile decrease with increasing					
porous parameter K _p . The temperature and concentration					
boosts up for distinct K values.					



Figure 11. Concentration on *diverse* of K_p.

	C _{fx}	Nu	Sh
M = 1.0	-1.9251	2.0193	0.3694
M = 1.5	-2.1237	1.9094	0.3640
M = 2.0	-2.2972	1.8105	0.3598
M = 2.5	-2.4518	1.7204	0.3566
$\beta = 1.0$	-1.3188	2.1424	0.3608
$\beta = 1.5$	-1.1767	2.1376	0.3534
$\beta = 2.0$	-1.1009	2.1334	0.3492
$\beta = 2.5$	-1.0536	2.1300	0.3464
$K_{p} = 1.0$	-2.1617	1.9066	0.3644
$K_{p} = 1.5$	-2.3353	1.8073	0.3602
$K_{p} = 2.0$	-2.4896	1.7169	0.3569
$K_{p} = 2.5$	-2.6289	1.6338	0.3542



Figure 12. Angular velocity on diverse of Kp.

The fallout of friction drag, Nusselt and Sherwood numeral for distinct of the momentum, casson parameter and porosity are shown in table. The Nusselt and Sherwood numeral retarded by increment of the momentum, casson parameter and porosity (M, β , Kp). Conversely, while Skin friction rises by gain of casson parameter whereas drops by M,Kp.

4.0 Conclusions

In this research numerical solution for the outcome of micro polar casson nano fluid with Porous medium are explored. MATLAB Bvp4c is used to solve the ODEs. The impact of planted variables on the disposal of mass, heat, porous medium and velocity functions is depicted graphically. The major findings are

- The Magnetic parameter, Casson parameter and Porous medium parameter are increased in temperature and concentration profile.
- Table 1 granted Nusselt and Sherwood numbers reduce by increasing the magnetic effect.
- The Velocity profile and Angular velocity profile decreases with enhancement in M, and K_p.

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Nomenclature

(<i>u</i> , <i>v</i>)	Velocity components	K_{1}	Permeability of the fluid	14	TT (1
β_{τ}	Thermal coefficient	q'''	Irregular heat parameter	M	Hartmann number
ß	Concentration coefficient	D	Mass diffusivity	N_{r}	Non-linear radiative parameter
	Fluid temperature	m C	Concentration susceptibility	S _r	Soret number
T T	Uniform tomporature	V V	Thermal diffusion ratio	S _C	Schmidt number
				D_u	Dufour number
C		1 _m	Mean temperature	E_{c}	Eckert number
$C_{_{\infty}}$	Uniform concentration	B_{iT}	Thermal Biot Number	ĸ	Chemical parameter
α	Angle of inclination	B_{ic}	Mass Biot Number	r a	Radiative heat
σ	Electric conductivity fluid	r	Velocity slip parameter	q_r K	Non dimensional parameter
B_{o}	Uniform magnetic field	G_{rT}	Thermal Grashot number	R D	
g	Gravitational acceleration	G_{rc}	Solutal Grashot number	P_r	Prandu number